

Precision predictions for hadron collider physics

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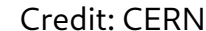


Figure 1 consists of nine Feynman diagrams labeled (a) through (i).
 (a) A quark q and antiquark \bar{q} annihilate into two gluons g .
 (b) Two gluons g annihilate into two gluons g .
 (c) Two gluons g annihilate into two gluons g .
 (d) A particle X^+ and an antiparticle X^- annihilate into a photon γ .
 (e) A particle X^+ and an antiparticle X^- annihilate into a Z boson.
 (f) A particle X^+ and an antiparticle X^- annihilate into a Higgs boson H .
 (g) A neutrino ν_e and an antineutrino $\bar{\nu}_e$ annihilate into a W boson.
 (h) A quark q and antiquark \bar{q} annihilate into a W boson.
 (i) A quark q and antiquark \bar{q} annihilate into a W boson and a photon γ or Z boson.

Credit: Jack Lindon, CERN

ATLAS Preliminary
 $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$

Events / GeV

Legend:
 ◆ Data
 — Fit
 - - - Background

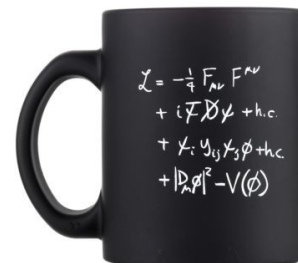
$H \rightarrow \gamma\gamma, m_H = 125.09 \text{ GeV}$

Data-Background

$m_{\gamma\gamma} [\text{GeV}]$

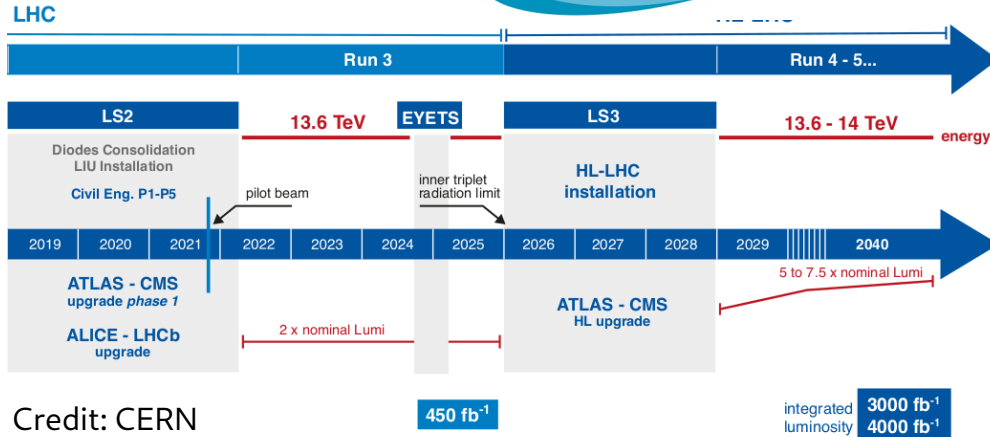
Credit: ATLAS

Credit: ATLAS



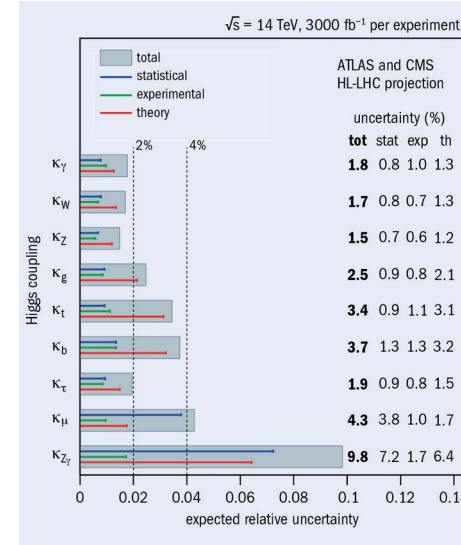
Credit: CERN

LHC Precision era and future experiments



Credit: CERN

Example:
projected Higgs couplings measurements



[1902.00134]

Theory input needed:

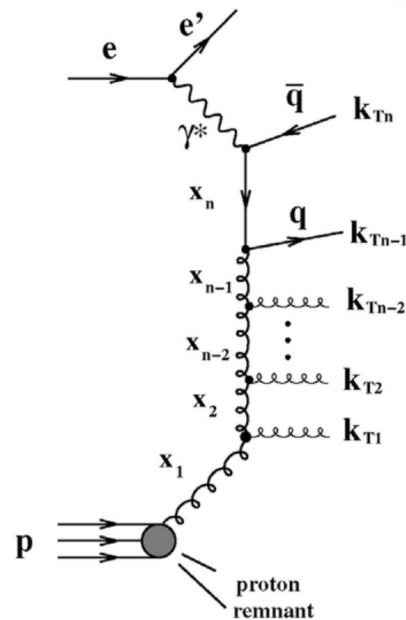
- Accurate → avoid wrong interpretation of excesses
- Precise → getting most out of our precious experiments

A pillar of precision phenomenology at hadron colliders

[Dokshitzer–Gribov–Lipatov–Altarelli–Parisi, '77]

DGLAP:

$$Q^2 \frac{\partial}{\partial Q^2} \begin{pmatrix} q_i(x, Q^2) \\ \bar{q}_i(x, Q^2) \\ g(x, Q^2) \end{pmatrix} = \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{d\xi}{\xi} \begin{pmatrix} P_{q_i q_j}(x/\xi) & 0 & P_{q_i g}(x/\xi) \\ 0 & P_{\bar{q}_i \bar{q}_j}(x/\xi) & P_{\bar{q}_i g}(x/\xi) \\ P_{g q_j}(x/\xi) & P_{g \bar{q}_j}(x/\xi) & P_{g g}(x/\xi) \end{pmatrix} \begin{pmatrix} q_j(\xi, Q^2) \\ \bar{q}_j(\xi, Q^2) \\ g(\xi, Q^2) \end{pmatrix}$$



Evolution of the proton PDF with the energy:

- Resummation of large logarithms from collinear emissions
- correlation of processes at different energies
- allows precise PDF determination

Basis of precision computation
at hadron colliders today

[Zeus, hep-ph/0502029]

Precision through higher orders

Hadronic cross section:

$$\sigma_{h_1 h_2 \rightarrow X} = \sum_{ij} \int_0^1 \int_0^1 dx_1 dx_2 \underbrace{\phi_{i,h_1}(x_1, \mu_F^2)}_{\text{PDFs}} \underbrace{\phi_{j/h_2}(x_2, \mu_F^2)}_{\text{DGLAP}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(\alpha_s(\mu_R^2), \mu_R^2, \mu_F^2)}_{\text{DGLAP}}$$

PDFs → DGLAP

Perturbative expansion of partonic cross section:

$$\hat{\sigma}_{ab \rightarrow X} = \underbrace{\alpha_s^0 \hat{\sigma}_{ab \rightarrow X}^{(0)}}_{\text{Leading order}} + \underbrace{\alpha_s^1 \hat{\sigma}_{ab \rightarrow X}^{(1)}}_{\text{Next-to-leading order}} + \underbrace{\alpha_s^2 \hat{\sigma}_{ab \rightarrow X}^{(2)}}_{\text{Next-to-next-to-leading order}} + \mathcal{O}(\alpha_s^3)$$

Leading order

Next-to-leading order

Next-to-next-to-leading order

Uncertainty:
 $\alpha_s(m_Z) \approx 0.118$

Order of magnitude

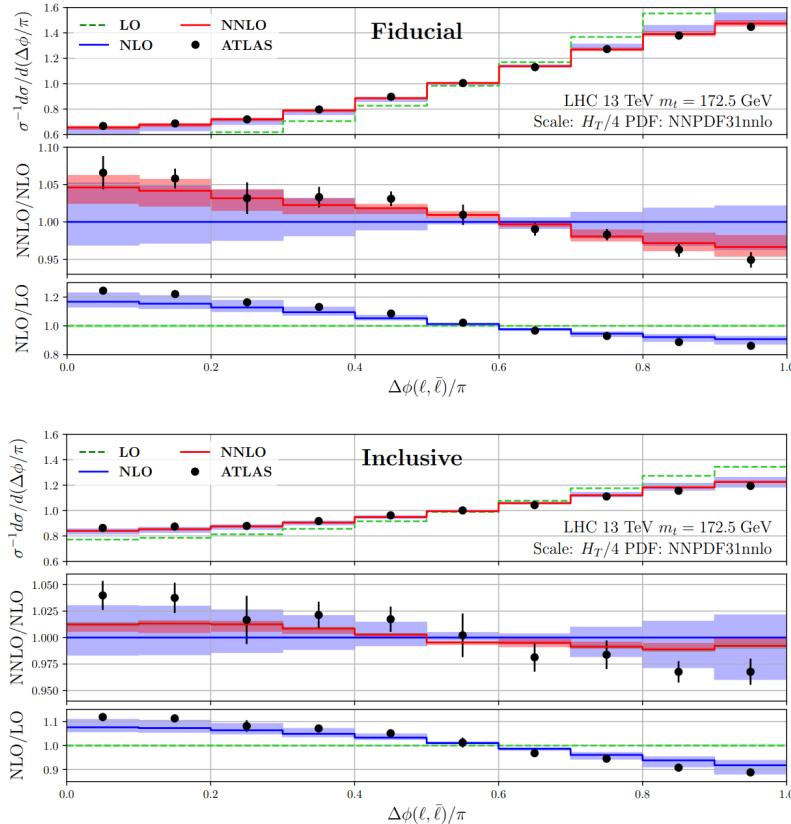
O(10%)

O(1%)

Next-to-next-to-leading order QCD needed to match experimental precision!
→ in some cases even next-to-next-to-next-to-leading order!

Accurate predictions, example: spin-correlations

Azimuthal correlations for leptons



[Behring, Czakon, Mitov, Papanastasiou, Poncelet'19
Czakon, Mitov, Poncelet '21]

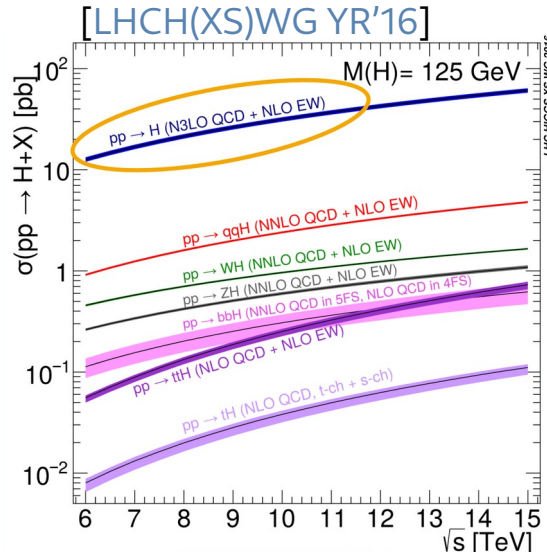
Spin-density-matrix

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_1^i d \cos \theta_2^j} = \frac{1}{4} \left(1 + B_1^i \cos \theta_1^i + B_2^j \cos \theta_2^j - C_{ij} \cos \theta_1^i \cos \theta_2^j \right)$$

Coefficient	LO ($\times 10^3$)	NLO ($\times 10^3$)	NNLO ($\times 10^3$)	CMS ($\times 10^3$)
B_1^k	1_{-0}^{+0} [sc] ± 1 [mc]	1_{-1}^{+0} [sc] ± 2 [mc]	-1_{-1}^{+0} [sc] ± 4 [mc]	5 ± 23
B_1^r	0_{-0}^{+0} [sc] ± 1 [mc]	0_{-0}^{+1} [sc] ± 2 [mc]	0_{-2}^{+1} [sc] ± 2 [mc]	-23 ± 17
B_1^n	0_{-0}^{+0} [sc] ± 1 [mc]	3_{-1}^{+1} [sc] ± 1 [mc]	4_{-0}^{+1} [sc] ± 3 [mc]	6 ± 13
B_2^k	0_{-0}^{+0} [sc] ± 1 [mc]	0_{-1}^{+0} [sc] ± 1 [mc]	-5_{-3}^{+2} [sc] ± 3 [mc]	7 ± 23
B_2^r	0_{-0}^{+0} [sc] ± 1 [mc]	0_{-0}^{+2} [sc] ± 1 [mc]	-2_{-1}^{+0} [sc] ± 2 [mc]	-10 ± 20
B_2^n	0_{-0}^{+0} [sc] ± 1 [mc]	-2_{-1}^{+0} [sc] ± 1 [mc]	-3_{-0}^{+1} [sc] ± 3 [mc]	17 ± 13
C_{kk}	324_{-7}^{+7} [sc] ± 1 [mc]	330_{-2}^{+2} [sc] ± 3 [mc]	323_{-5}^{+2} [sc] ± 6 [mc]	300 ± 38
C_{rr}	6_{-5}^{+5} [sc] ± 1 [mc]	58_{-12}^{+18} [sc] ± 2 [mc]	69_{-7}^{+8} [sc] ± 3 [mc]	81 ± 32
C_{nn}	332_{-0}^{+1} [sc] ± 1 [mc]	330_{-1}^{+1} [sc] ± 2 [mc]	326_{-1}^{+1} [sc] ± 4 [mc]	329 ± 20
$C_{nr} + C_{rn}$	1_{-0}^{+0} [sc] ± 1 [mc]	-1_{-0}^{+1} [sc] ± 3 [mc]	-4_{-0}^{+4} [sc] ± 6 [mc]	-4 ± 37
$C_{nr} - C_{rn}$	0_{-1}^{+0} [sc] ± 1 [mc]	-1_{-0}^{+1} [sc] ± 2 [mc]	2_{-2}^{+4} [sc] ± 8 [mc]	-1 ± 38
$C_{nk} + C_{kn}$	0_{-0}^{+0} [sc] ± 1 [mc]	2_{-0}^{+1} [sc] ± 1 [mc]	3_{-1}^{+4} [sc] ± 3 [mc]	-43 ± 41
$C_{nk} - C_{kn}$	1_{-0}^{+0} [sc] ± 1 [mc]	1_{-1}^{+1} [sc] ± 2 [mc]	6_{-0}^{+2} [sc] ± 7 [mc]	40 ± 29
$C_{rk} + C_{kr}$	-229_{-4}^{+4} [sc] ± 1 [mc]	-203_{-7}^{+9} [sc] ± 2 [mc]	-194_{-6}^{+8} [sc] ± 7 [mc]	-193 ± 64
$C_{rk} - C_{kr}$	1_{-0}^{+0} [sc] ± 1 [mc]	1_{-1}^{+0} [sc] ± 4 [mc]	-1_{-3}^{+1} [sc] ± 5 [mc]	57 ± 46

[CMS 1907.03729]

Precision example: Quark-mass effects in Higgs production

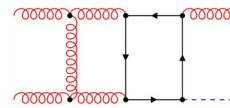
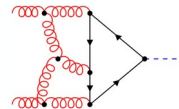
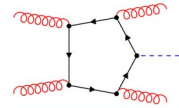
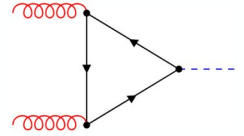


Second-order corrections to top-bottom interference effects with full mass dependence!

[Czakon, Eschment, Niggetiedt, Poncelet, Schellenberger: 2312.09896, 2407.12413]

Higgs-production in gluon fusion, main uncertainties:

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
+0.10 pb -1.15 pb	± 0.18 pb	± 0.56 pb	± 0.49 pb	± 0.40 pb	± 0.49 pb
+0.21% -2.37%	$\pm 0.37\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$



↓

Renorm. scheme	$\overline{\text{MS}}$	on-shell
$\mathcal{O}(\alpha_s^2)$	-1.11	-1.98
LO	$-1.11^{+0.28}_{-0.43}$	$-1.98^{+0.38}_{-0.53}$
$\mathcal{O}(\alpha_s^3)$	-0.65	-0.44
NLO	$-1.76^{+0.27}_{-0.28}$	$-2.42^{+0.19}_{-0.12}$
$\mathcal{O}(\alpha_s^4)$	+0.02	+0.43
NNLO	$-1.74(2)^{+0.13}_{-0.03}$	$-1.99(2)^{+0.29}_{-0.15}$

$\overline{\text{MS}}$ vs. on-shell scheme:
→ first agreement at NNLO!

Precision example: strong-coupling from TEEC

ATLAS

Particle-level TEEC

$\sqrt{s} = 13 \text{ TeV}; 139 \text{ fb}^{-1}$

anti- k_t $R = 0.4$

$p_T > 60 \text{ GeV}$

$|\eta| < 2.4$

$\mu_{R,F} = \hat{p}_T$

$\alpha_s(m_Z) = 0.1180$

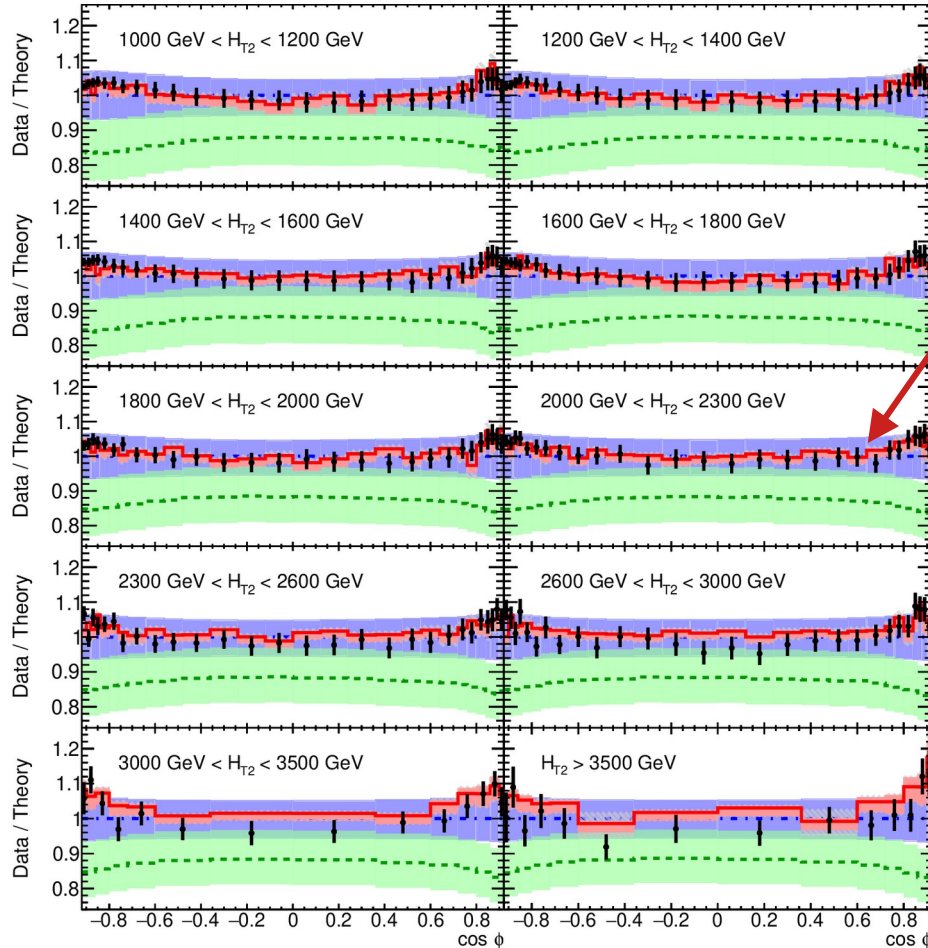
NNPDF 3.0 (NNLO)

— Data

--- LO

--- NLO

--- NNLO



NNLO QCD for three jets

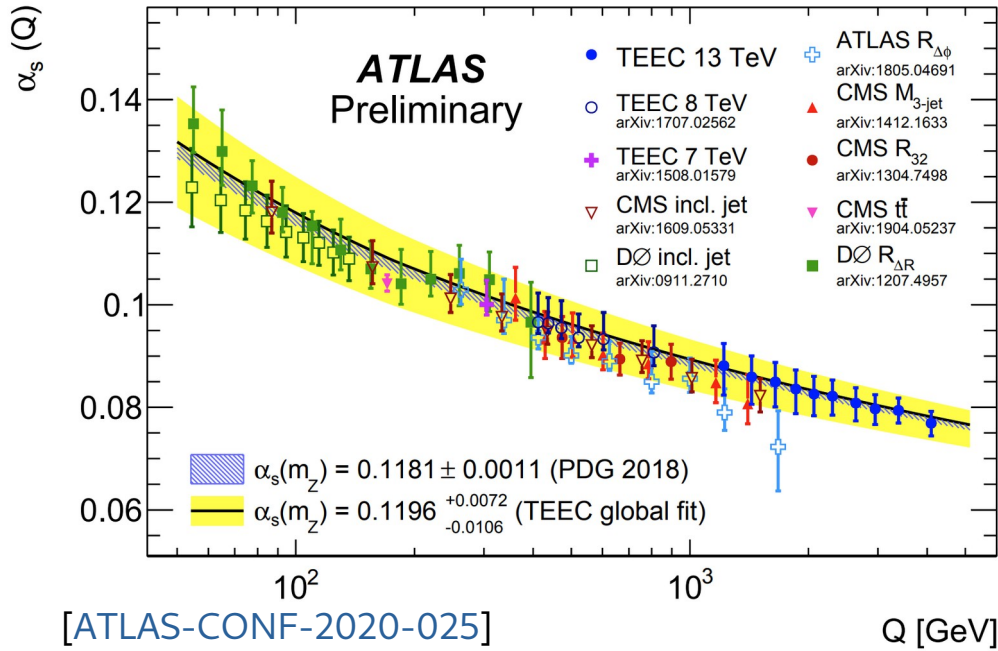
- [Czakon, Mitov, Poncelet 2106.05331]
- [Alvarez, Cantero, Czakon, Llorente, Mitov, Poncelet 2301.01086]

ATLAS: α_s extraction

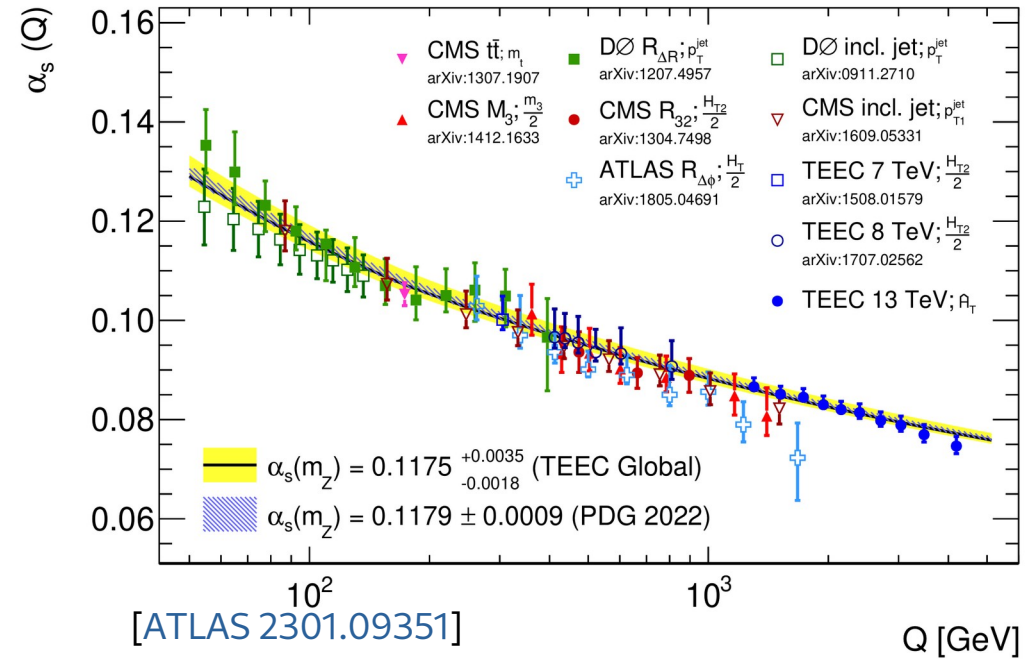
- [ATLAS 2301.09351]

Precision example: strong-coupling from TEEC

NLO QCD



NNLO QCD



From calculations to phenomenology

Resummation/
Fragmentation

Multi-loop
amplitudes

Higher-order
computations

Numerical
integration

Subtraction
schemes

Precision
phenomenology

Impact on data interpretation!

Conclusions

Predictions are essential for data interpretation. We need them

- accurate → avoid wrong interpretation of excesses
- precise → getting most out of our precious experiments

QCD calculations went a long way from '77 to today:

- Miles stones like NNLO QCD multi-jet production and N3LO QCD for simple inclusive processes
- Next challenge: full automation of NNLO QCD and incorporation into Monte Carlos (main bottlenecks: multi-loop amplitudes, parton-shower matching)
- A lot of space for surprises and novel ideas!

**Many thanks to all my collaborators,
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Thank you!**