

Polarization modelling in MBI processes/ Precision Predictions for Polarized Electroweak Bosons

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LEVERHULME
TRUST



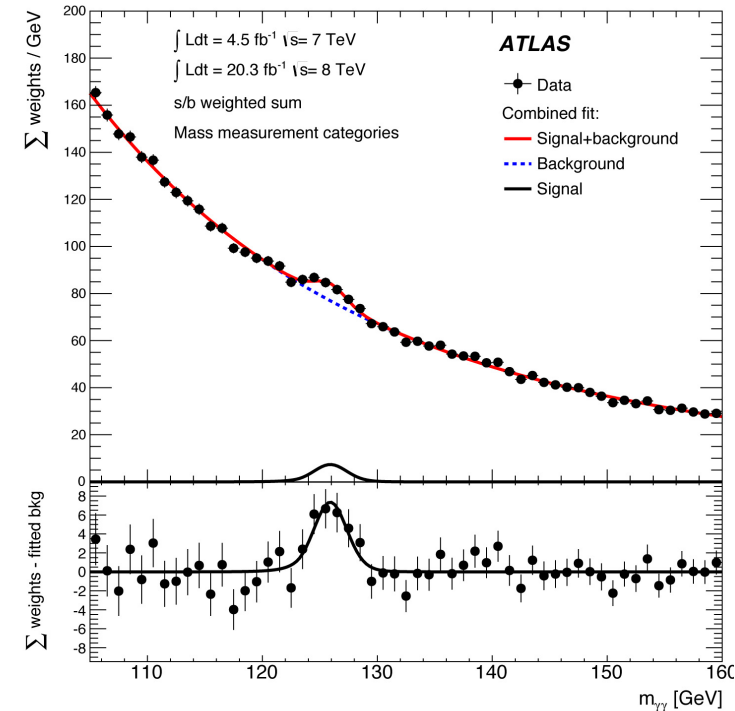
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Motivation: Why polarized boson?

- 10 Years Jubilee of the discovery of a “Higgs”-like boson
- New era of precision measurements of Higgs properties: spin, couplings, self-interactions, ...
To answer fundamental questions like
 - Is the Higgs the Standard Model Higgs?
 - Is it a fundamental scalar?
 - How is the EWSB mechanism realized?
 - ...
- **Polarized bosons** can be a tool to access these questions from a point of view **orthogonal** to Higgs measurements.



Motivation: Why polarized boson?

The reason is the EWSB in the SM:

$$\mathcal{L}_{\text{EW}} = -\frac{1}{4}(W_{\mu\nu}^i)^2 - \frac{1}{4}(B_{\mu\nu}^i)^2 + (D_\mu\phi)^2 - V(\phi^\dagger\phi)$$

- Higgs potential and minimum:

$$V(\phi^\dagger\phi) = -\mu^2(\phi^\dagger\phi) + \lambda(\phi^\dagger\phi)^2 \quad \phi = U(\pi^i) \begin{pmatrix} 0 \\ \frac{v+H}{\sqrt{2}} \end{pmatrix} \quad \text{Vev: } \phi^\dagger\phi = \frac{\mu^2}{2\lambda} \equiv \frac{v^2}{2}$$

- Goldstone bosons can be absorbed via gauge transformation (unitary gauge).
This gives rise to massive gauge bosons:

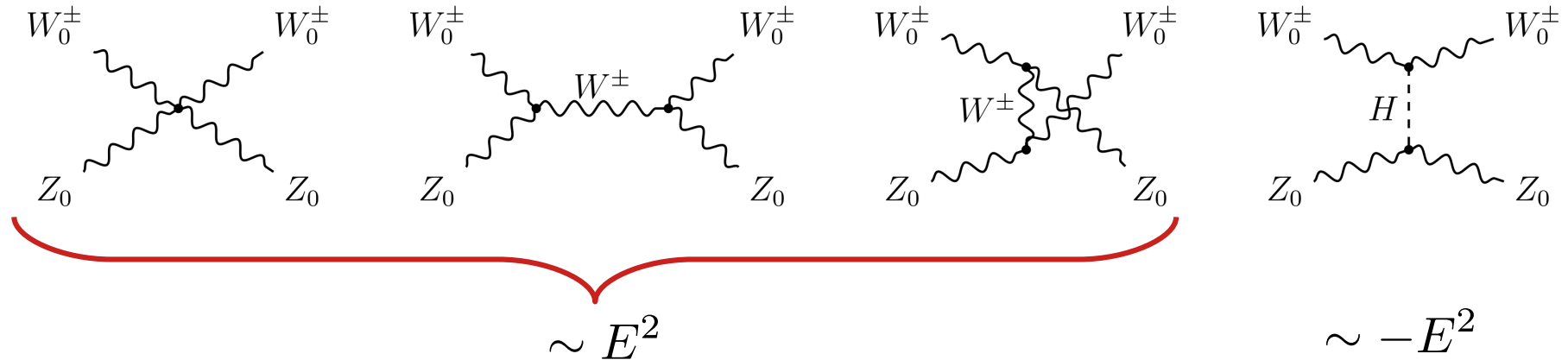
$$\phi = U^{-1}(\pi^i)\phi, \quad W_\mu = U^{-1}W_\mu U - \frac{i}{g_W}U^{-1}\partial_\mu U$$

$$|D_\mu\phi|^2 \ni \frac{v^2}{8} [2g_W^2 W_\mu^+ W^{-\mu} + (g_W W_\mu^3 - g'_W B_\mu)^2] \quad \longrightarrow \quad M_W = \frac{1}{2}vg_W, \quad M_Z = \frac{M_W}{\cos\theta_W}$$

- Restores renormalizability and unitarity

Motivation: Why polarized boson?

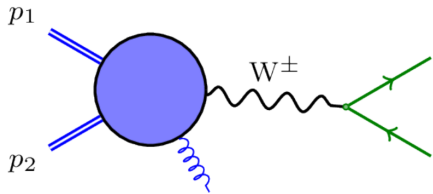
Scattering of longitudinal vector bosons:



- Measurement of polarized boson scattering or production probes:
 - EWSB mechanism
 - Higgs and gauge sector
 - New physics models (for example throw EFTs like SMEFT)

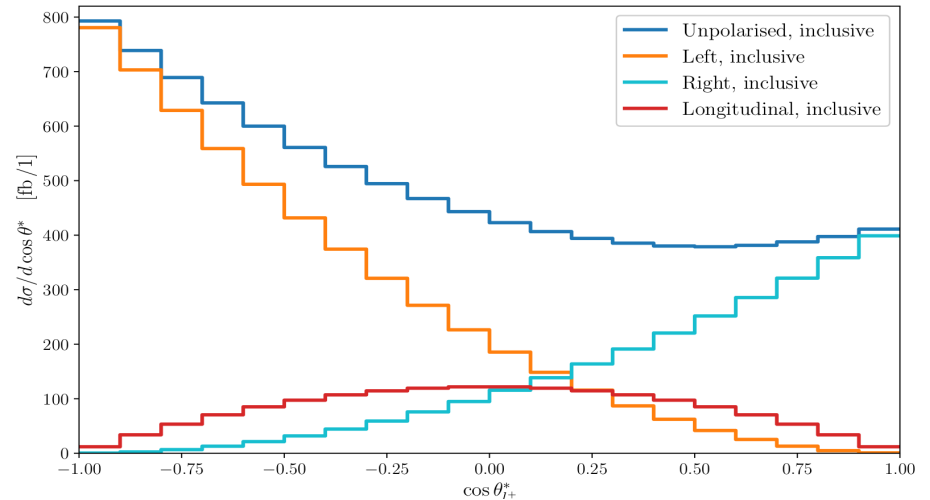
How to measure polarized bosons?

- No WW collider... but the LHC produces massive vector bosons abundantly!
- We can't measure boson polarization directly.
- Luckily decay products can be used as a "polarimeter":



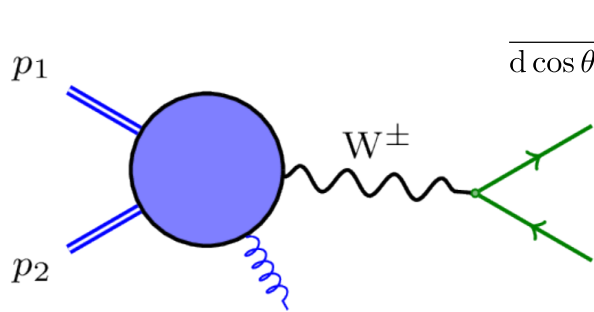
$$M_\lambda = \mathbf{P}_\mu \cdot \frac{-g_{\mu\nu} + \frac{k^\mu k^\nu}{k^2}}{k^2 - M_V^2 + iM_V\Gamma_V} \cdot \mathbf{D}_\nu$$

$$\left(-g^{\mu\nu} + \frac{k^\mu k^\nu}{k^2}\right) \rightarrow \sum_\lambda \epsilon_\lambda^{*\mu} \epsilon_\lambda^\nu$$



How to measure polarized bosons

Angular decomposition of W decay:



$$\frac{d\sigma}{d\cos\theta^* d\phi^* dX} = \frac{d\sigma}{dX} \frac{3}{16\pi} \left[(1 + \cos^2\theta^*) + \frac{A_0}{2}(1 - 3\cos^2\theta^*) + A_1 \sin 2\theta^* \cos\phi^* + \frac{A_2}{2} \sin^2\theta^* \cos 2\phi^* \right. \\ \left. + A_3 \sin\theta^* \cos\phi^* + A_4 \cos\theta^* + A_5 \sin^2\theta^* \sin 2\phi^* + A_6 \sin 2\theta^* \sin\phi^* + A_7 \sin\theta^* \sin\phi^* \right]$$

After azimuthal integration:

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta^*} = \frac{3}{4} \sin^2\theta^* f_0 + \frac{3}{8} (1 - \cos\theta^*)^2 f_L + \frac{3}{8} (1 + \cos\theta^*)^2 f_R$$

Idea: Suitable projections (or fits) extract fractions of left, right and longitudinal components.

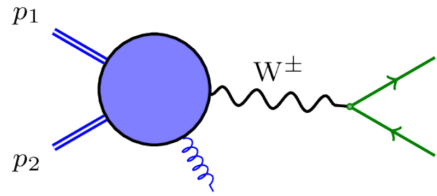
Reality check

This simple idea suffers from:

- Fiducial phase space requirements on the leptons
 - Interferences do not cancel \rightarrow nice correspondence between fractions (f_0, f_R, f_L) and angular distributions broken.
- Higher order corrections to decay
 - Decomposition in $\{A_i\}$ does not hold any more
- Angles in boson rest frame
 - Z rest frame accessible, but W more difficult to reconstruct

The more general solution is to generate polarized events!

Polarized cross sections



$$M_\lambda = \mathbf{P}_\mu \cdot \frac{-g_{\mu\nu} + \frac{k^\mu k^\nu}{k^2}}{k^2 - M_V^2 + iM_V\Gamma_V} \cdot \mathbf{D}_\nu$$

On-shell bosons: (DPA or NWA) $\left(-g^{\mu\nu} + \frac{k^\mu k^\nu}{k^2}\right) \rightarrow \sum_\lambda \epsilon_\lambda^{*\mu} \epsilon_\lambda^\nu$

$$|M|^2 = \underbrace{\sum_\lambda |M_\lambda|^2}_{\text{pol. x-section}} + \underbrace{\sum_{\lambda \neq \lambda'} M_\lambda^* M_{\lambda'}}_{\text{Interferences}}$$

→ pol. x-section Interferences

Create samples of specified cross section

$$\frac{d\sigma}{dX} = f_L \frac{d\sigma_L}{dX} + f_R \frac{d\sigma_R}{dX} + f_0 \frac{d\sigma_0}{dX} \left(+ f_{int.} \frac{d\sigma_{int.}}{dX} \right)$$

And fit f_L, f_R, f_0 to measured $\frac{d\sigma^{exp.}}{dX}$

Advantages:

- Interferences can be handled
- X can be any observable → lab frame observables
- $\frac{d\sigma_i}{dX}$ can be systematically improved

Overview SM results

- **Polarised VBS (so far LO):**

W boson polarization in vector boson scattering at the LHC,

Ballestrero, Maina, Pelliccioli 1710.09339

Polarized vector boson scattering in the fully leptonic WZ and ZZ channels at the LHC,

Ballestrero, Maina, Pelliccioli 1907.04722

- **Automated predictions from polarized matrix elements**

- Buarque Franzosi, Mattelaer, Ruiz, Shil 1912.01725

Different polarization definitions in same-sign WW scattering at the LHC,

Ballestrero, Maina, Pelliccioli 2007.07133

- **Polarized Diboson NLO QCD / NLO EW : WW / WZ / ZZ**

Fiducial polarization observables in hadronic WZ production: A next-to-leading order QCD+EW study,

Baglio, Le Duc 1810.11034

Anomalous triple gauge boson couplings in ZZ production at the LHC and the role of Z boson polarizations,

Rahama, Singh 1810.11657

Polarization observables in WZ production at the 13 TeV LHC: Inclusive case,

Baglio, Le Duc 1910.13746

Unravelling the anomalous gauge boson couplings in ZW⁺- production at the LHC and the role of spin-1 polarizations,

Rahama, Singh 1911.03111

Polarized electroweak bosons in W+W- production at the LHC including NLO QCD effects,

Denner, Pelliccioli 2006.14867

NLO EW and QCD corrections to polarized ZZ production in the four-charged-lepton channel at the LHC,

Denner, Pelliccioli 2107.06579

Breaking down the entire spectrum of spin correlations of a pair of particles involving fermions and gauge bosons,

Rahama, Singh 2109.09345

NLO QCD predictions for doubly-polarized WZ production at the LHC,

Denner, Pelliccioli 2010.07149

Doubly-polarized WZ hadronic cross sections at NLO QCD+EW accuracy,

Duc Ninh Le, Baglio 2203.01470

Doubly-polarized WZ hadronic production at NLO QCD+EW: Calculation method and further results

Duc Ninh Le, Baglio, Dao 2208.09232

Overview SM results

- Polarized Diboson NNLO QCD

NNLO QCD study of polarised W+W- production at the LHC,
Poncelet, Popescu 2102.13583

- Single boson production

Left-Handed W Bosons at the LHC,

Z. Bern et. al. 1103.5445

Electroweak gauge boson polarisation at the LHC,

Stirling, Vryonidou 1204.6427

What Does the CMS Measurement of W-polarization Tell Us about the Underlying Theory of the Coupling of W-Bosons to Matter?,

Belyaev, Ross 1303.3297

Polarised W+j production at the LHC: a study at NNLO QCD accuracy,

Pellen, Poncelet, Popescu 2109.14336

- Experimental works on polarized W production

Measurement of the Polarization of W Bosons with Large Transverse Momenta in W+Jets Events at the LHC,

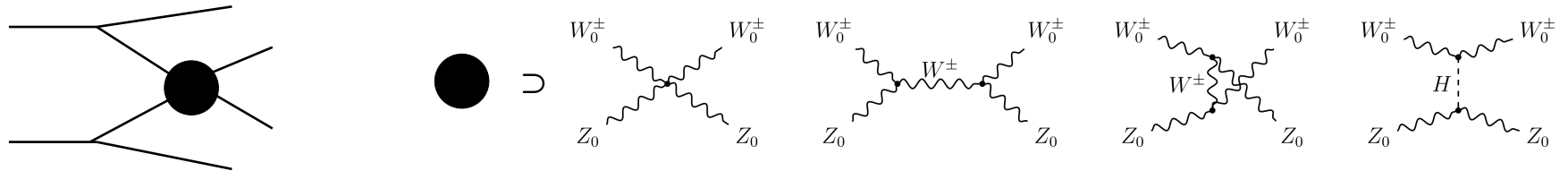
CMS 1104.3829

Measurement of the polarisation of W bosons produced with large transverse momentum in pp collisions at $\sqrt{s}=7$ TeV with the ATLAS experiment,

ATLAS 1203.2165

Polarized VBS

Clear connection to unitarity cancellation:



Extensive studies within the PHANTOM MC and Madgraph at LO:

- Singly and doubly polarized cross sections
- Frame of polarization definition
- Impact of fiducial selection criteria
- Study of off-shell/interference effects

W boson polarization in vector boson scattering at the LHC,

Ballestrero, Maina, Pelliccioli 1710.09339

Polarized vector boson scattering in the fully leptonic WZ and ZZ channels at the LHC,

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Automated predictions from polarized matrix elements

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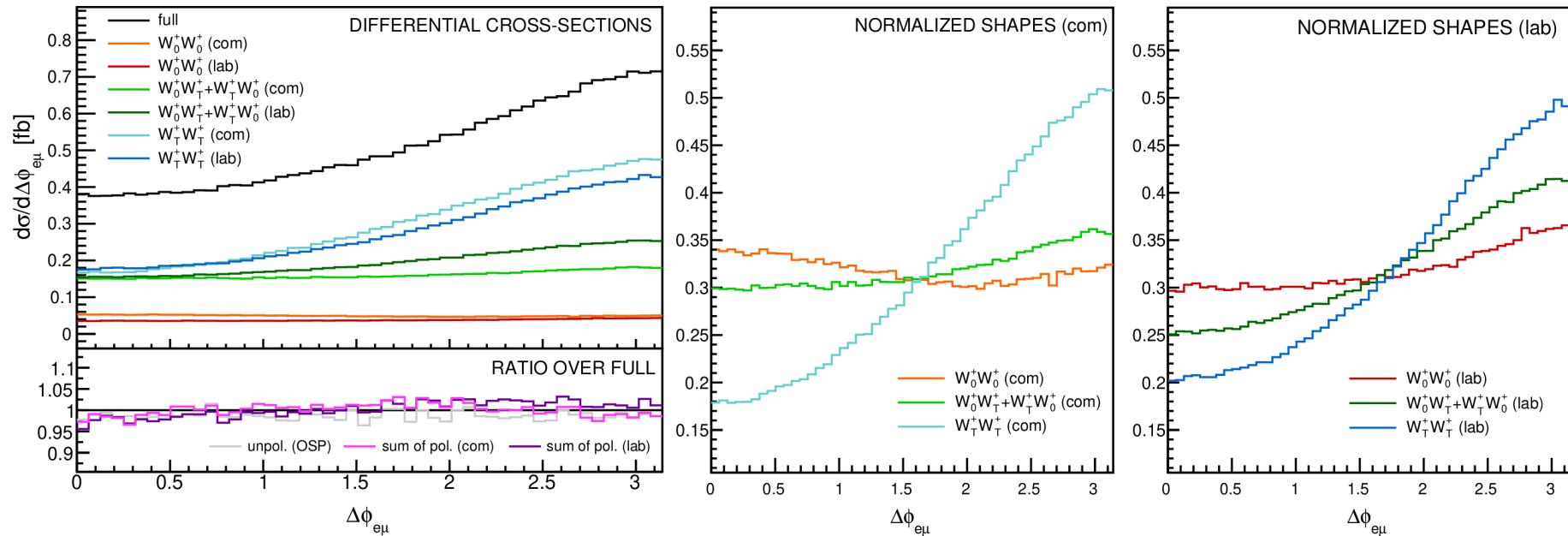
Different polarization definitions in same-sign WW scattering at the LHC,

Ballestrero, Maina, Pelliccioli 2007.07133

Polarized VBS: same-sign WW

Different polarization definitions in same-sign WW scattering at the LHC,
 Ballestrero, Maina, Pelliccioli 2007.07133

- Visibly different shapes for double longitudinal polarized (LL) component
- However small total contribution of LL: $\sim 10\%$ of total cross section.



Polarized VBS: same-sign WW

Different polarization definitions in same-sign WW scattering at the LHC,
Ballestrero, Maina, Pelliccioli 2007.07133

- Visibly different shapes for double longitudinal polarized (LL) component
- However small total contribution of LL: $\sim 10\%$ of total cross section.
=> Similar also for other VBS channels

W+W+

	Lab	WW CoM	ratio
full	3.185(3)	-	-
unpol	3.167(2)	-	-
0-unpol	0.8772(8)	0.8374(9)	0.95
T-unpol	2.287(2)	2.329(2)	1.02
0-0	0.2573(3)	0.3275(4)	1.27
0-T, T-0	0.6199(6)	0.5081(5)	0.82
T-T	1.666(1)	1.820(1)	1.09

W+W-

	Lab	WW CoM	ratio
full	4.651(2)	-	-
unpol	4.641(2)	-	-
0-unpol	1.186(1)	1.146(1)	0.97
T-unpol	3.456(2)	3.494(2)	1.01
unpol-0	1.2226(4)	1.1905(5)	0.97
unpol-T	3.418(1)	3.450(1)	1.01
0-0	0.3314(2)	0.3786(3)	1.14
0-T	0.8545(4)	0.7669(3)	0.90
T-0	0.8912(4)	0.8119(4)	0.91
T-T	2.563(1)	2.683(1)	1.05

WZ

	Lab	WZ CoM	ratio
full	0.5253(3)	-	-
unpol	0.5210(3)	-	-
0-unpol	0.1216(1)	0.1292(1)	1.06
T-unpol	0.3992(2)	0.3918(3)	0.98
unpol-0	0.1370(1)	0.1436(1)	1.05
unpol-T	0.3839(2)	0.3773(2)	0.98
0-0	0.03236(3)	0.03993(5)	1.23
0-T	0.08923(8)	0.08926(8)	1.00
T-0	0.1045(1)	0.1039(1)	0.99
T-T	0.2948(2)	0.2876(2)	0.98

pp \rightarrow ZZ \rightarrow 4l NLO EW/QCD

NLO EW and QCD corrections to polarized ZZ production in the four-charged-lepton channel at the LHC,
Denner, Pelliccioli 2107.06579

$$pp \rightarrow ZZ \rightarrow 4l + X$$

Integrated cross sections:

- Small LL contribution
- Sizeable QCD and EW corrections
- Large gg-loop induced contribution
- Fractions preserved from LO to NLO

mode	σ_{LO} [fb]	δ_{QCD}	δ_{EW}	δ_{gg}	σ_{NLO^+} [fb]	$\sigma_{\text{NLO}^\times}$ [fb]
full	11.1143(5) ^{+5.6%} _{-6.8%}	+34.9%	-11.0%	+15.6%	15.505(6) ^{+5.7%} _{-4.4%}	15.076(5) ^{+5.5%} _{-4.2%}
unpol.	11.0214(5) ^{+5.6%} _{-6.8%}	+35.0%	-10.9%	+15.7%	15.416(5) ^{+5.7%} _{-4.4%}	14.997(4) ^{+5.5%} _{-4.2%}
Z _L Z _L	0.64302(5) ^{+6.8%} _{-8.1%}	+35.7%	-10.2%	+14.5%	0.9002(6) ^{+5.5%} _{-4.3%}	0.8769(5) ^{+5.4%} _{-4.1%}
Z _L Z _T	1.30468(9) ^{+6.5%} _{-7.7%}	+45.3%	-9.9%	+2.8%	1.8016(9) ^{+4.3%} _{-3.5%}	1.7426(8) ^{+4.1%} _{-3.3%}
Z _T Z _L	1.30854(9) ^{+6.5%} _{-7.7%}	+44.3%	-9.9%	+2.8%	1.7933(9) ^{+4.3%} _{-3.4%}	1.7355(8) ^{+4.0%} _{-3.2%}
Z _T Z _T	7.6425(3) ^{+5.2%} _{-6.4%}	+31.2%	-11.2%	+20.5%	10.739(4) ^{+6.2%} _{-4.7%}	10.471(3) ^{+6.1%} _{-4.6%}

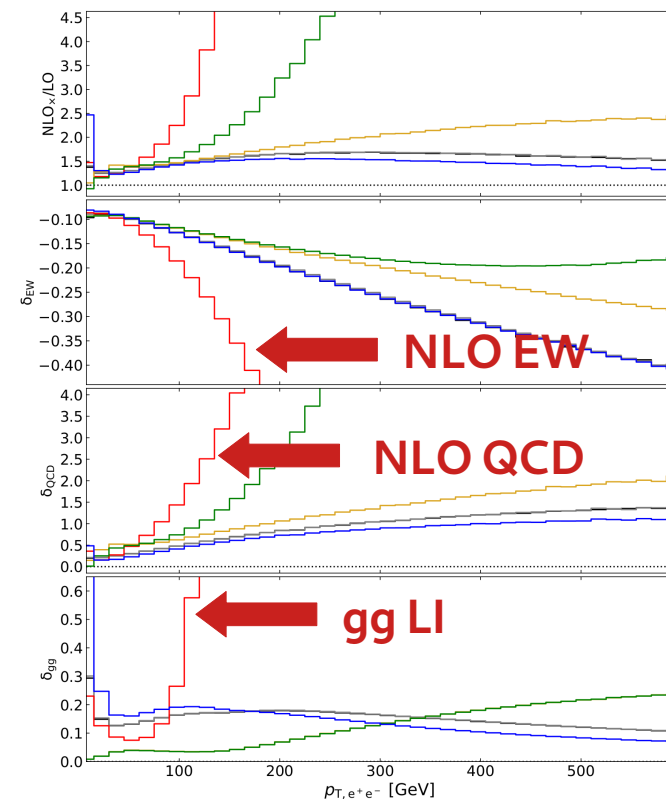
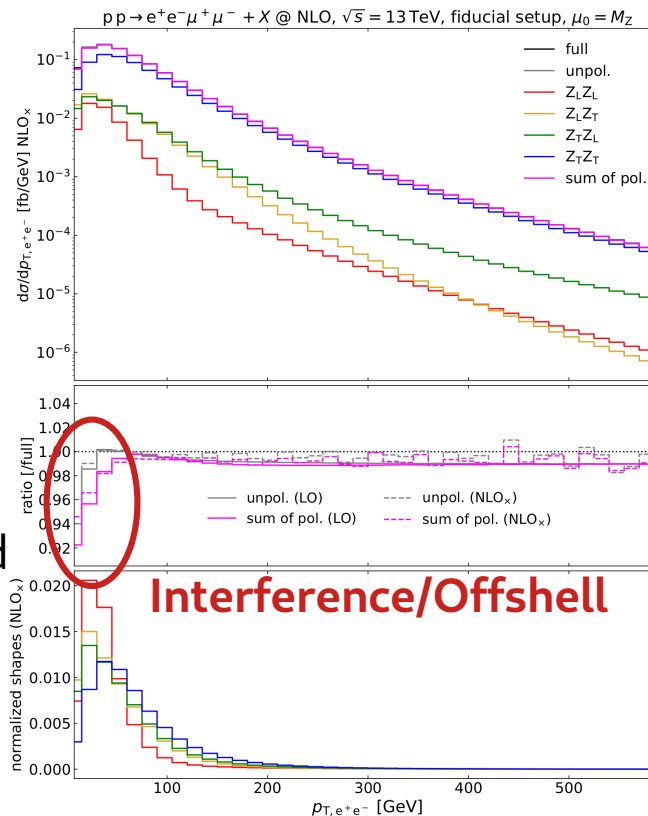
pp \rightarrow ZZ \rightarrow 4l NLO EW/QCD

NLO EW and QCD corrections to polarized ZZ production in the four-charged-lepton channel at the LHC,
Denner, Pelliccioli 2107.06579

Differential cross section

- Very large QCD, EW and gg-LI corrections for **LL contribution**
- Comparison to full predictions reveals interferences and off-shell effects
- Similar findings for WZ and WW final states, for example in:

Doubly-polarized WZ hadronic cross sections at NLO QCD+EW accuracy,
Duc Ninh Le, Baglio 2203.01470



pp \rightarrow W+W- \rightarrow e+ mu- ve vm NNLO QCD

NNLO QCD study of polarised W+W- production at the LHC,
Poncelet, Popescu 2102.13583

LL gets largest NNLO
QCD corrections

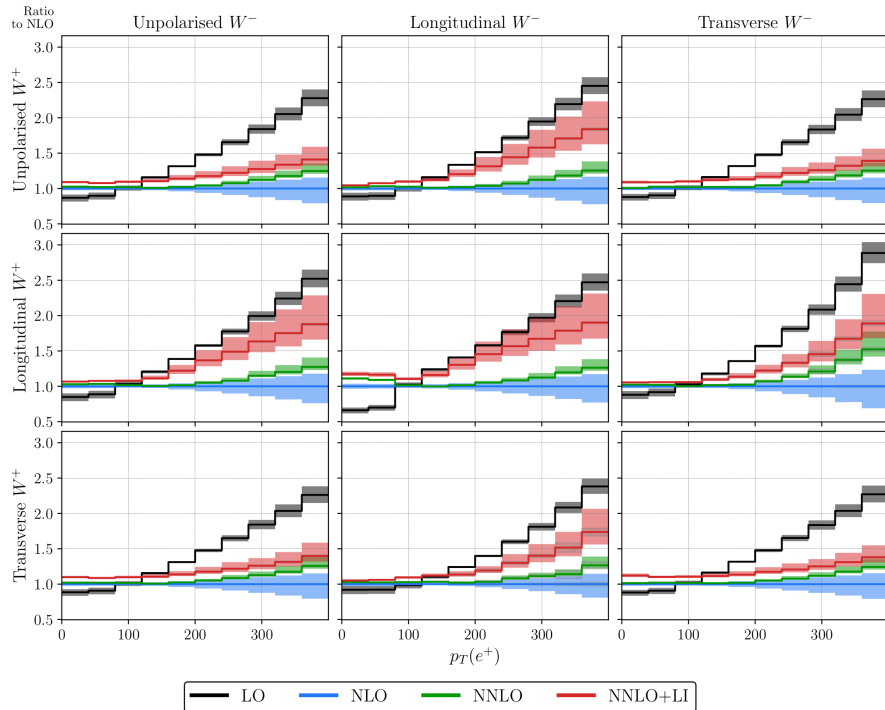
LI leads to sizeable
corrections and
enhanced scale
dependence

	NLO	NNLO	K_{NNLO}	LI	NNLO+LI
off-shell	220.06(5) ^{+1.8%} _{-2.3%}	225.4(4) ^{+0.6%} _{-0.6%}	1.024	13.8(2) ^{+25.5%} _{-18.7%}	239.1(4) ^{+1.5%} _{-1.2%}
unpol. (nwa)	221.85(8) ^{+1.8%} _{-2.3%}	227.3(6) ^{+0.6%} _{-0.6%}	1.025	13.68(3) ^{+25.5%} _{-18.7%}	241.0(6) ^{+1.5%} _{-1.1%}
unpol. (dpa)	214.55(7) ^{+1.8%} _{-2.3%}	219.4(4) ^{+0.6%} _{-0.6%}	1.023	13.28(3) ^{+25.5%} _{-18.7%}	232.7(4) ^{+1.4%} _{-1.1%}
W_L^+ (dpa)	57.48(3) ^{+1.9%} _{-2.6%}	59.3(2) ^{+0.7%} _{-0.7%}	1.032	2.478(6) ^{+25.5%} _{-18.3%}	61.8(2) ^{+1.0%} _{-0.8%}
W_L^- (dpa)	63.69(5) ^{+1.9%} _{-2.6%}	65.4(3) ^{+0.8%} _{-0.8%}	1.026	2.488(6) ^{+25.5%} _{-18.3%}	67.9(3) ^{+0.9%} _{-0.8%}
W_T^+ (dpa)	152.58(9) ^{+1.7%} _{-2.1%}	155.7(6) ^{+0.7%} _{-0.6%}	1.020	11.19(2) ^{+25.5%} _{-18.8%}	166.9(6) ^{+1.6%} _{-1.3%}
W_T^- (dpa)	156.41(7) ^{+1.7%} _{-2.1%}	159.7(6) ^{+0.5%} _{-0.6%}	1.021	11.19(2) ^{+25.5%} _{-18.8%}	170.9(6) ^{+1.7%} _{-1.3%}
$W_L^+ W_L^-$ (dpa)	9.064(6) ^{+3.0%} _{-3.0%}	9.88(3) ^{+1.3%} _{-1.3%}	1.090	0.695(2) ^{+25.5%} _{-18.8%}	10.57(3) ^{+2.9%} _{-2.4%}
$W_L^+ W_T^-$ (dpa)	48.34(3) ^{+1.9%} _{-2.5%}	49.4(2) ^{+0.9%} _{-0.7%}	1.021	1.790(5) ^{+25.5%} _{-18.3%}	51.2(2) ^{+0.6%} _{-0.8%}
$W_T^+ W_L^-$ (dpa)	54.11(5) ^{+1.9%} _{-2.5%}	55.5(4) ^{+0.6%} _{-0.7%}	1.025	1.774(5) ^{+25.5%} _{-18.3%}	57.2(4) ^{+0.7%} _{-0.7%}
$W_T^+ W_T^-$ (dpa)	106.26(4) ^{+1.6%} _{-1.9%}	108.3(3) ^{+0.5%} _{-0.5%}	1.019	9.58(2) ^{+25.5%} _{-18.9%}	117.9(3) ^{+2.1%} _{-1.6%}

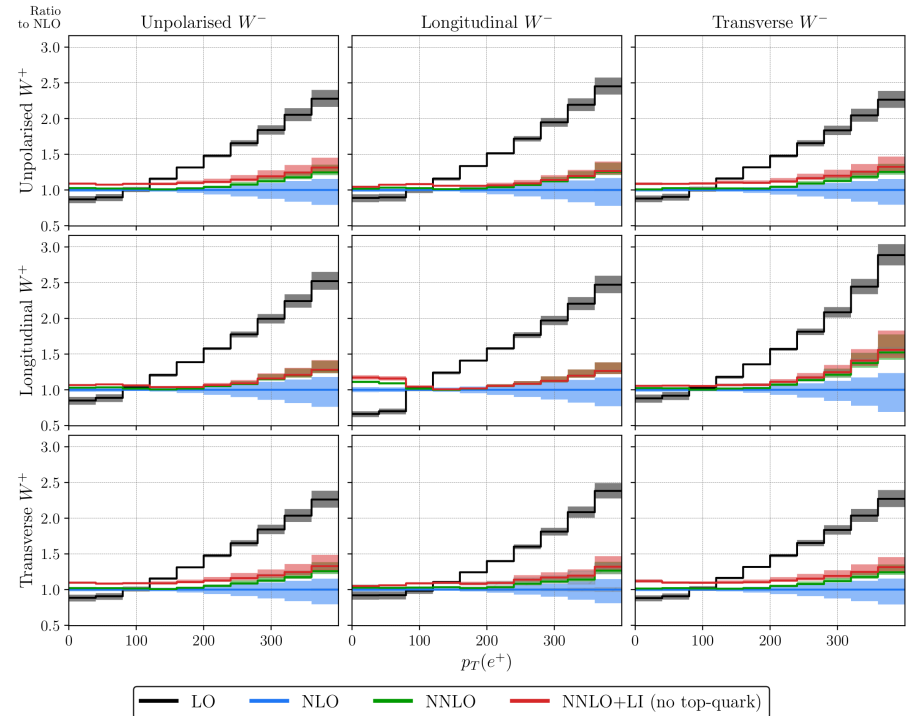
pp \rightarrow W+W- \rightarrow e+ mu- ve vm NNLO QCD

NNLO QCD study of polarised W+W- production at the LHC,
Poncelet, Popescu 2102.13583

With top-quark loops in gg LI



Without top-quark loops in gg LI



Polarized W+jet @ NNLO QCD

Polarised W+j production at the LHC: a study at NNLO QCD accuracy, Pellen, Poncelet, Popescu 2109.14336

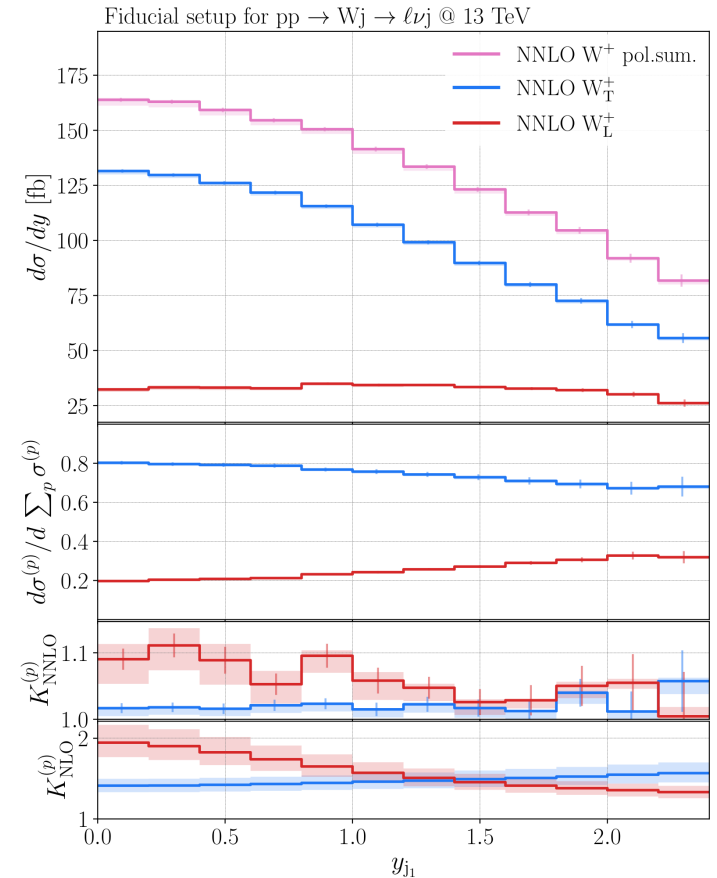
NNLO QCD predictions for polarized W production in fiducial and “inclusive” phase space

- Similar findings as in diboson: longitudinal component gets large QCD corrections:
 - up to 100% at NLO and 10 % NNLO
 - significant reduction of scale dependence
- How do higher order corrections impact extractions of polarization fractions?

$$\frac{d\sigma}{dX} = f_T \frac{d\sigma_T}{dX} + f_L \frac{d\sigma_L}{dX} \left(+ \cancel{f_{int} \frac{d\sigma_{int.}}{dX}} \right)$$

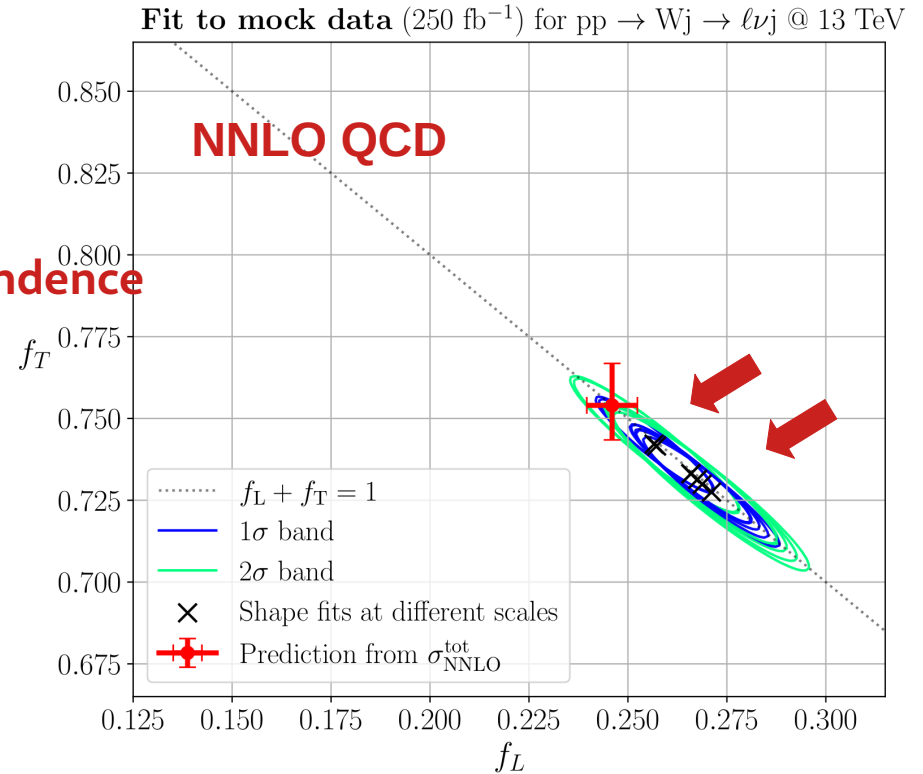
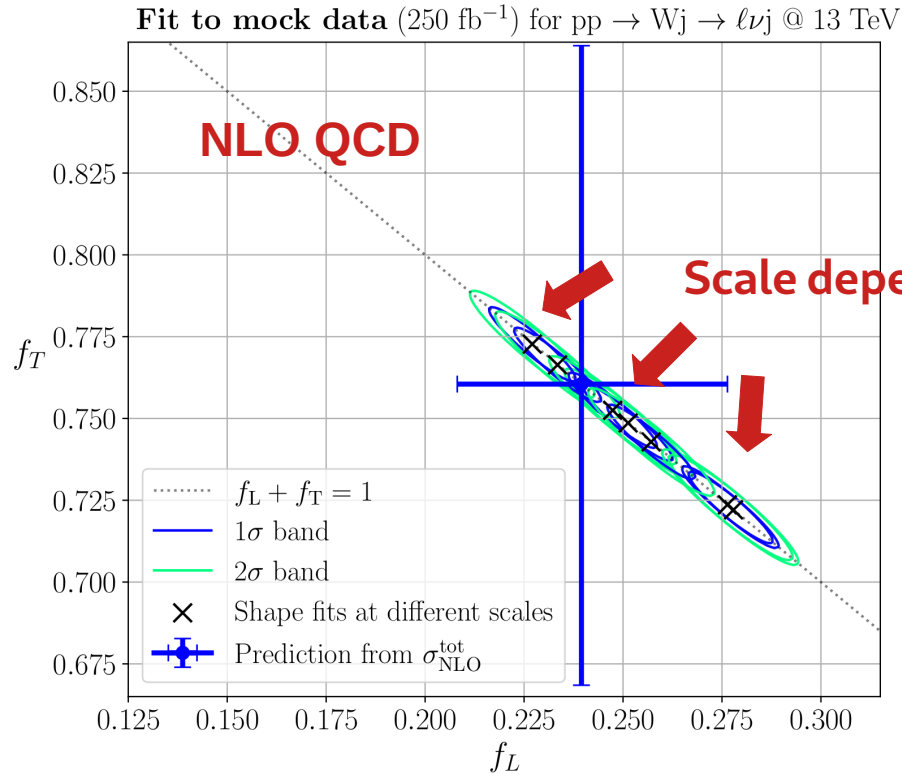
→ 0

(by selecting regions with little interferences)



W+jet: mock-data fit

First: Fit to mock-data (based on NNLO QCD and 250 fb⁻¹ stats):
→ extreme case to see effect of scale dependence reduction

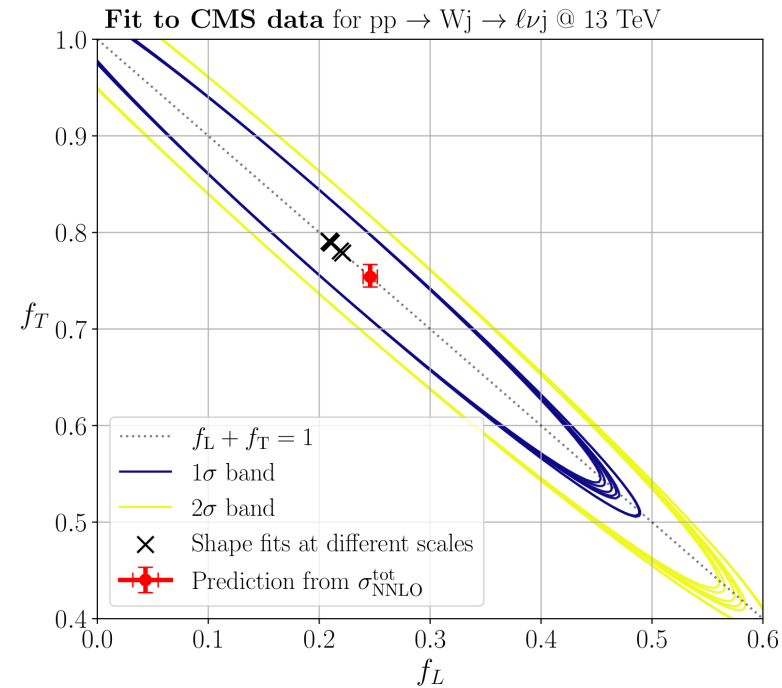
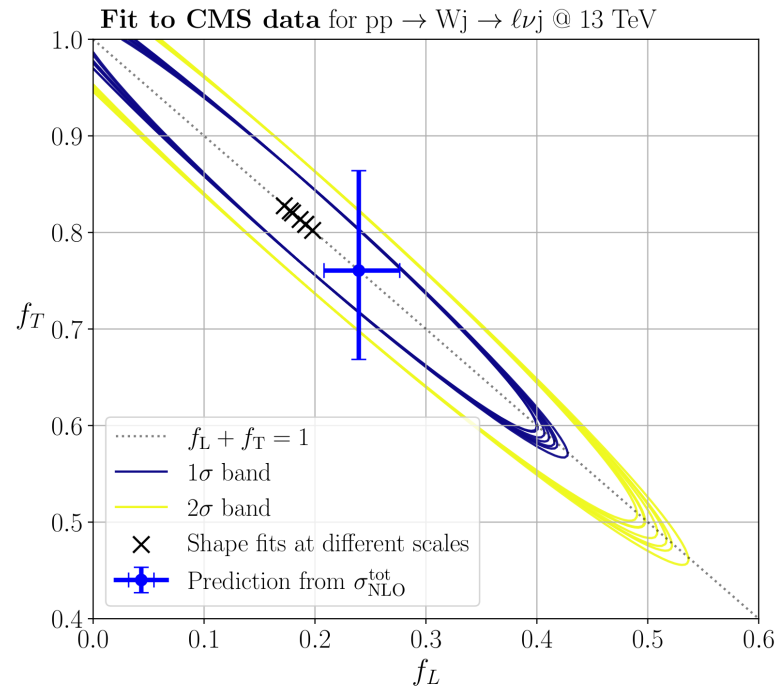


W+jet : fit to CMS data

Second: Fit to actual data

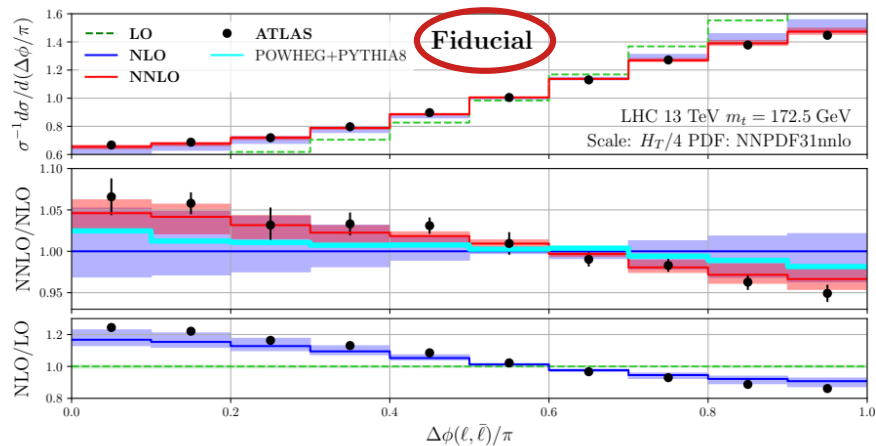
Measurement of the differential cross sections for the associated production of a W boson and jets in proton-proton collisions at $\sqrt{s}=13$ TeV, CMS 1707.05979

→ so far dominated by experimental uncertainties:

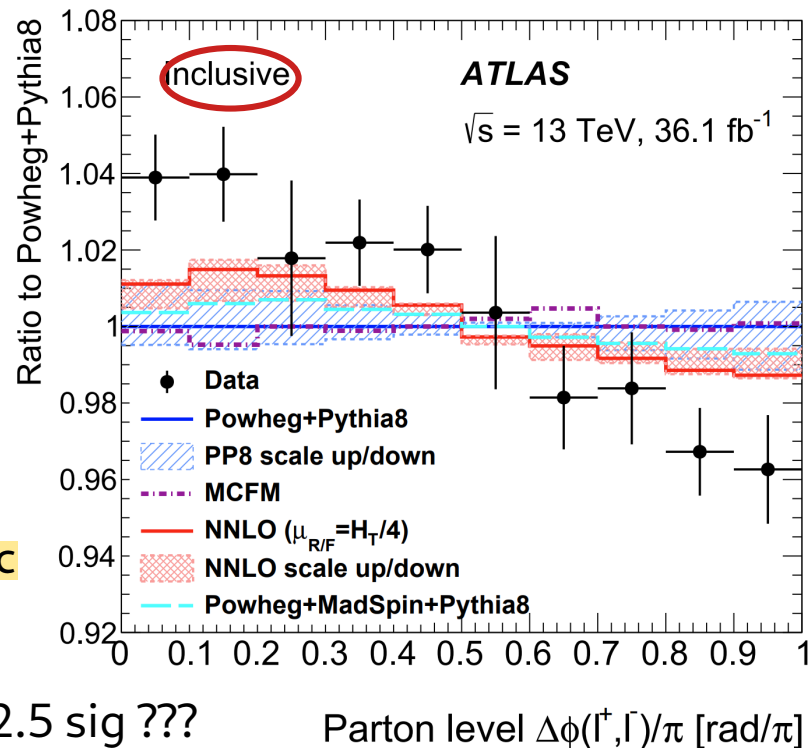


Remark: Lessons from ttbar

Measurements of top-quark pair spin correlations in the e mu channel at $\sqrt{s} = 13$ TeV using pp collisions in the ATLAS detector, ATLAS 1903.07570



Higher order corrections to spin correlations in top quark pair production at the LHC
Behring, Czakon, Mitov, Papanastasiou, Poncelet 1901.05407



Extrapolation with
Powheg+Pythia (NLO+PS) \rightarrow ~ 2.5 sig ???
Higher order corrections! (in red)

Parton level $\Delta\phi(\Gamma^+, \Gamma^-)/\pi$ [rad/ π]

Conclusion & outlook

Summary:

- Increasing interest in studying polarized bosons
→ triggered by exciting prospects for future precise measurements
- Higher order corrections are crucial to measure polarization fractions accurately.
→ Efforts to provide fixed order predictions at (N)NLO QCD and NLO EW
Diboson and single boson final states: WW , WZ , ZZ , W +jet

Outlook:

- More realistic simulations require parton shower effects.
- Higher order corrections for single-boson or boson pairs
→ Corrections to polarized VBS?

Thank you!